

ON THE ROOTS OF THE LEGENDRE, LAGUERRE, AND HERMITE POLYNOMIALS

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ABSTRACT. *For several orthogonal polynomials, Cohen proved that their roots are the eigenvalues of symmetric tridiagonal matrices. In this paper, we give examples of this Cohen’s result for the Legendre, Laguerre, and Hermite polynomials, which are useful in applications to quantum mechanics and numerical analysis.*

KEYWORDS: Laguerre and Hermite polynomials, Leverrier-Takeno’s technique, Legendre polynomials.

INTRODUCTION

Here we consider the Legendre polynomials $P_n(x)$ [1, 2]:

$$P_0(x) = 1, \quad P_1(x) = x, \quad P_2(x) = \frac{1}{2}(3x^2 - 1), \quad P_3(x) = \frac{1}{2}(5x^3 - 3x), \dots \quad (1)$$

which verify the differential equation $(1 - x^2)y'' - 2xy' + l(l + 1) = 0$, $l = 0, 1, 2, \dots$; Cohen [3, 4] showed that roots of $P_n(x) = 0$ are the proper values of the following symmetric tridiagonal matrix:

$$P_n = \begin{pmatrix} 0 & \frac{1}{\sqrt{3}} & \square & \square & \square & \square & \square \\ \frac{1}{\sqrt{3}} & 0 & \frac{2}{\sqrt{15}} & \square & \square & \square & \square \\ \square & \frac{2}{\sqrt{15}} & 0 & \frac{3}{\sqrt{35}} & \square & \square & \square \\ \square & \square & \frac{3}{\sqrt{35}} & 0 & \ddots & \square & \square \\ \square & \square & \square & \ddots & \ddots & \ddots & \square \\ \square & \square & \square & \square & \ddots & 0 & g(n) \\ \square & \square & \square & \square & \square & g(n) & 0 \end{pmatrix}, \quad g(n) = \frac{n-1}{\sqrt{(2n-1)(2n-3)}}. \quad (2)$$

Since the eigenvalues of a symmetric matrix are all real, it follows that the roots of the Legendre polynomials must all be real [1, 5, 6]. Moreover, the absence of nonzero terms along the leading diagonal of the matrix P_n implies that the eigenvalues are symmetrically

SOME APPLICATIONS OF THE COHEN'S RESULTS

The characteristic equation of a matrix $A_{n \times n}$:

$$\lambda^n + a_1 \lambda^{n-1} + a_2 \lambda^{n-2} + \dots + a_n = 0, \quad (7)$$

can be constructed via the Leverrier-Takeno's procedure:

$$a_1 = -s_1, \quad a_2 = \frac{1}{2}[(s_1)^2 - s_2], \quad a_3 = \frac{1}{6}[-(s_1)^3 + 3s_1 s_2 - 2s_3], \dots \quad (8)$$

where s_r is the trace of A^r . Then we consider (2):

$$P_2 = \begin{pmatrix} 0 & \frac{1}{\sqrt{3}} \\ \frac{1}{\sqrt{3}} & 0 \end{pmatrix}, \quad P_2^2 = \begin{pmatrix} \frac{1}{3} & 0 \\ 0 & \frac{1}{3} \end{pmatrix}, \quad s_1 = a_1 = 0, \quad s_2 = \frac{2}{3}, \quad a_2 = -\frac{1}{3},$$

thus (7) implies the equation $3\lambda^2 - 1 = 0$ in agreement with $P_2(x) = 0$. Similarly:

$$P_3 = \begin{pmatrix} 0 & \frac{1}{\sqrt{3}} & 0 \\ \frac{1}{\sqrt{3}} & 0 & \frac{2}{\sqrt{15}} \\ 0 & \frac{2}{\sqrt{15}} & 0 \end{pmatrix}, \quad P_3^2 = \begin{pmatrix} \frac{1}{3} & 0 & \frac{2}{3\sqrt{5}} \\ 0 & \frac{3}{5} & 0 \\ \frac{2}{3\sqrt{5}} & 0 & \frac{4}{15} \end{pmatrix}, \quad P_3^3 = \begin{pmatrix} 0 & \frac{\sqrt{3}}{5} & 0 \\ \frac{\sqrt{3}}{5} & 0 & \frac{6}{5\sqrt{15}} \\ 0 & \frac{6}{5\sqrt{15}} & 0 \end{pmatrix},$$

Hence, $s_1 = a_1 = 0$, $s_2 = \frac{6}{5}$, $a_2 = -\frac{3}{5}$, $s_3 = a_3 = 0$, and from (7) we obtain $5\lambda^3 - 3\lambda = 0$ in harmony with $P_3(x) = 0$. Let's remember that the roots of Legendre polynomials are important in the Gaussian quadrature [7], in the study of electromagnetic radiation and the angular function for the hydrogen atom.

Besides, from (4):

$$L_2 = \begin{pmatrix} 1 & 1 \\ 1 & 3 \end{pmatrix}, \quad L_2^2 = \begin{pmatrix} 2 & 4 \\ 4 & 10 \end{pmatrix}, \quad s_1 = 4, \quad s_2 = 12, \quad a_1 = -4, \quad a_2 = 2,$$

then (7) gives $\lambda^2 - 4\lambda + 2 = 0$, equivalent to $L_2(x) = 0$; and:

$$L_3 = \begin{pmatrix} 1 & 1 & 0 \\ 1 & 3 & 2 \\ 0 & 2 & 5 \end{pmatrix}, \quad L_3^2 = \begin{pmatrix} 2 & 4 & 2 \\ 4 & 14 & 16 \\ 2 & 6 & 29 \end{pmatrix}, \quad L_3^3 = \begin{pmatrix} 6 & 18 & 18 \\ 18 & 78 & 108 \\ 18 & 108 & 177 \end{pmatrix},$$

therefore $s_1 = 9$, $s_2 = 45$, $s_3 = 261$, $a_1 = -9$, $a_2 = 18$, $a_3 = -6$, and from (7) we deduce that $\lambda^3 - 9\lambda^2 + 18\lambda - 6 = 0$ in according with (3). The Laguerre polynomials participate in the radial function of hydrogen-like atoms [19] and diatomic molecules [20].

For the Hermite polynomials, we have:

$$H_2 = \begin{pmatrix} 0 & \sqrt{\frac{1}{2}} \\ \sqrt{\frac{1}{2}} & 0 \end{pmatrix}, \quad H_2^2 = \begin{pmatrix} \frac{1}{2} & 0 \\ 0 & \frac{1}{2} \end{pmatrix}, \quad s_1 = a_1 = 0, \quad s_2 = 1, \quad a_2 = -\frac{1}{2},$$

implying the characteristic equation $2\lambda^2 - 1 = 0$, equivalent to $H_2(x) = 0$; and:

$$H_3 = \begin{pmatrix} 0 & \sqrt{\frac{1}{2}} & 0 \\ \sqrt{\frac{1}{2}} & 0 & 1 \\ 0 & 1 & 0 \end{pmatrix}, \quad H_3^2 = \begin{pmatrix} \frac{1}{2} & 0 & \sqrt{\frac{1}{2}} \\ 0 & \frac{3}{2} & 0 \\ \sqrt{\frac{1}{2}} & 0 & 1 \end{pmatrix}, \quad H_3^3 = \begin{pmatrix} 0 & \frac{3}{2}\sqrt{\frac{1}{2}} & 0 \\ \frac{3}{2}\sqrt{\frac{1}{2}} & 0 & \frac{3}{2} \\ 0 & \frac{3}{2} & 0 \end{pmatrix},$$

$s_1 = a_1 = 0$, $s_2 = 3$, $a_2 = -\frac{3}{2}$, $s_3 = a_3 = 0$, thus (7) gives the expression $2\lambda^3 - 3\lambda = 0$ which is compatible with $H_3(x) = 0$. The Hermite polynomials are fundamental in the analysis of the harmonic oscillator in quantum physics.

Thus we have that the Leverrier-Takeno's process [18] allows to see that the eigenvalues of the matrices (2, 4, 6) are the roots of the Legendre, Laguerre [21, 22], and Hermite polynomials, respectively. Let's remember that the QR algorithm [23-26] is an efficient method to determine the proper values of a matrix.

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